

# Genuine CP-Odd Observables at the LHC and $ZZH$ Coupling

Tao Han \*

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MCTP Spring Symposium on Higgs Boson Physics

Ann Arbor, May 14, 2010

\*Collaborators: Neil Christensen and Y.-C. Li

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Concluding Remarks

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- Or, differential w.r.t. a *definitive* CP observable  $PS \rightarrow PS'$ :

$$dA_{CP} \propto d_{PS} \Gamma(i^0 \rightarrow f) - d_{PS'} \Gamma(i^0 \rightarrow \bar{f}),$$

*e.g.* :  $(p^+ \times p^-) \cdot p^0 \dots$

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With longitudinally polarized beams  $\sigma_1\sigma_2$ ,

$$e^-(p_1) e^+(p_2) \rightarrow \ell^-(q_1) \ell^+(q_2) X^0$$

under CP:

$$\mathcal{M}_{--}(\vec{p}_1, \vec{p}_2; \vec{q}_1, \vec{q}_2) \Rightarrow \mathcal{M}_{++}(\vec{p}_1, \vec{p}_2; -\vec{q}_2, -\vec{q}_1), \quad (1)$$

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Eq. (2), CP eigenstate: construct angular FB asymmetries:

$$d\mathcal{A}_{FB} = d\sigma_{-+}^F - d\sigma_{-+}^B.$$

Simple examples include:

$$\begin{aligned} e^-(p_1) e^+(p_2) &\rightarrow W^+ W^- \rightarrow l^-(q_1) l^+(q_2) \nu \bar{\nu}, \\ &\rightarrow \tau^+ \tau^- \rightarrow l^-(q_1) l^+(q_2) 2\nu 2\bar{\nu}, \\ &\rightarrow t\bar{t} \rightarrow l^-(q_1) l^+(q_2) \nu \bar{\nu} b\bar{b}, \\ &\rightarrow Z H^0 \rightarrow l^-(q_1) l^+(q_2) H^0, \\ &\rightarrow \chi^+ \chi^-, \chi_i^0 \chi_j^0, \tilde{l}^+ \tilde{l}^- \rightarrow l^-(q_1) l^+(q_2) X^0. \end{aligned}$$



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Such as CP-odd observables: ( $\vec{q}_\pm = \vec{q}_1 \pm \vec{q}_2$ )

$$\begin{array}{lll}
 \cos \theta_Z \sim \hat{z} \cdot \vec{q}_+, & \cos \theta_\ell \sim \hat{z} \cdot (\vec{q}_1 \times \vec{q}_2), & \sin \theta_- \sim (\hat{z} \times \vec{q}_-) \cdot (\vec{q}_1 \times \vec{q}_2). \\
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Note that:\*

<p>CP-odd, <math>\hat{T}</math>-even:  <math>\sin\phi_{CP} \sin\delta</math></p> <p>CP-even, <math>\hat{T}</math>-even:  <math>\cos\phi_{CP} \cos\delta</math></p>	<p>CP-odd, <math>\hat{T}</math>-odd:  <math>\sin\phi_{CP} \cos\delta</math></p> <p>CP-even, <math>\hat{T}</math>-odd:  <math>\cos\phi_{CP} \sin\delta</math></p>
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## Challenge: CP-odd observables at the LHC

Difficulties in choosing a state (process):

- $pp$  is NOT a CP eigenstate (comparing with  $\bar{p}\bar{p}$  ? )  
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(wish the absence of a CP-conserving phase  $\delta \approx 0$ .)

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(b). Or start from a well-defined  $i^0$ : §  
 $H^0(\text{or } X^0) \rightarrow ZZ \rightarrow \ell^+ \ell^- \ell^+ \ell^-.$

‡ Kane, Ladinsky, Yuan (1991); Dawson, Valencia (1997); Langacker, Paz, Wang, Yavin (2007); Kumar, Rajaraman, Wells (2008).

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- Transverse variables:

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or its variations like:

$$((\vec{p}_f \times \vec{p}_{\bar{f}}) \cdot \hat{z})^{2m+1} \cdot (\vec{p}_{fT} \cdot \vec{p}_{\bar{f}T})^n \operatorname{sgn}((\vec{p}_f - \vec{p}_{\bar{f}}) \cdot \hat{z}),$$

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## Example: General $ZZH$ Vertex

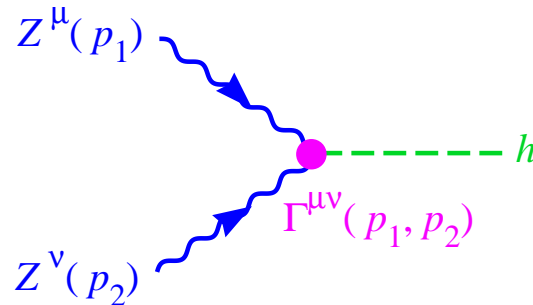
(for EWSB sector<sup>\*\*</sup>)

<sup>\*\*</sup>Yukawa  $t\bar{t}H$ : Gunion and He (1996); Godbole (2009).

## Example: General $ZZH$ Vertex

(for EWSB sector\*\*)

Most general vertex function for  $ZZh$



$$\Gamma^{\mu\nu}(p_1, p_2) = i \frac{2}{v} h [a M_Z^2 g^{\mu\nu} + b (p_1^\mu p_2^\nu - p_1 \cdot p_2 g^{\mu\nu}) + \tilde{b} \epsilon^{\mu\nu\rho\sigma} p_{1\rho} p_{2\sigma}]$$

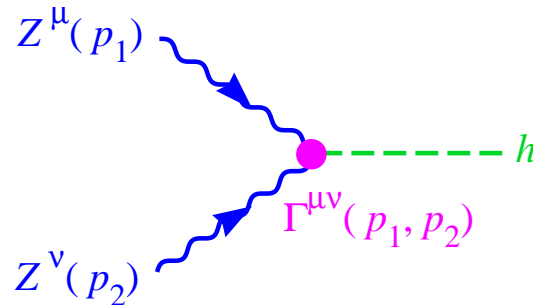
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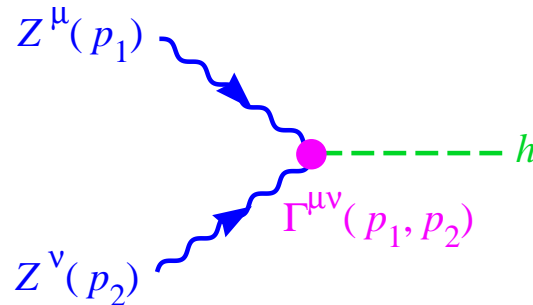
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In an “effective field theory”, operators  $\frac{g^2}{\Lambda^2} H H W^{\mu\nu} W_{\mu\nu}, \frac{g^2}{\Lambda^2} H H W^{\mu\nu} \tilde{W}_{\mu\nu}$ ,  
natural size:  $a, b, \tilde{b} \sim \mathcal{O}(\frac{1}{16\pi^2} \sim 1)$ .

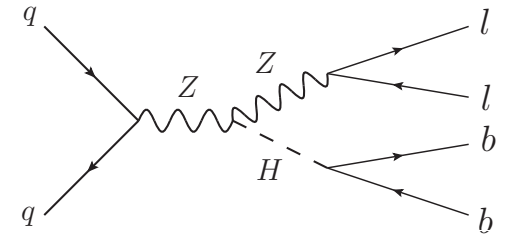
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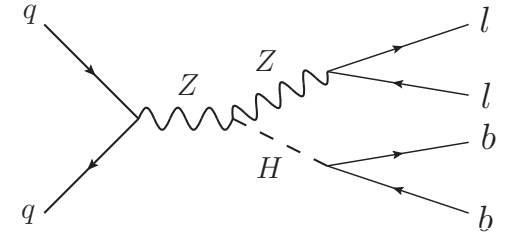


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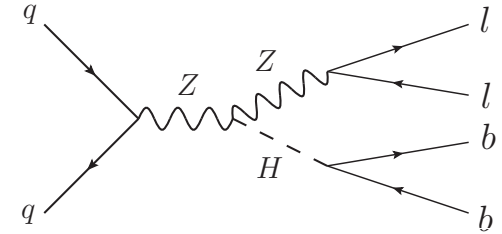
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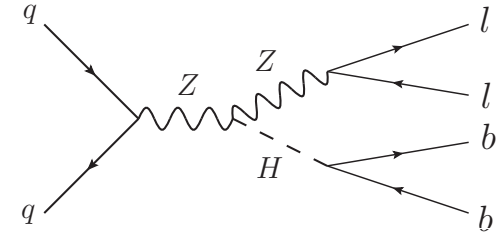
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- Vector-boson-fusion was exploited for  $ZZH$ ,  $WWH$ :<sup>‡‡</sup>  
 $qq' \rightarrow qq'H$ ,

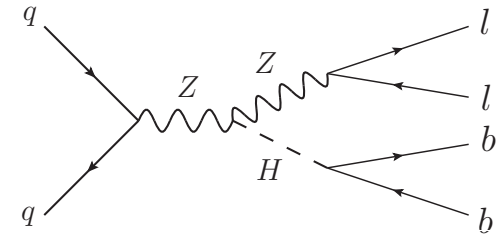
<sup>††</sup>Butterworth, Davison, Rubin, Salam (2008).

<sup>‡‡</sup>Plehn, Rainwater, Zeppenfeld (2002).

# Search at the LHC

$$pp \rightarrow \ell^- \ell^+ H X,$$

$$\text{or } q\bar{q} \rightarrow ZH \rightarrow \ell^- \ell^+ H(b\bar{b}).$$



- Appropriate channel for  $M_H < 140$  GeV, below  $ZZ$ .
- Recent studies improve the observability of this channel.<sup>††</sup>
- Vector-boson-fusion was exploited for  $ZZH$ ,  $WWH$ :<sup>‡‡</sup>  
 $qq' \rightarrow qq'H$ ,

with distinctive distributions for  $a, b$  and  $\tilde{b}$  terms,  
 but no charge info, no  $q$  versus  $\bar{q}$ ,  
 thus not a test of CP violation (actually only P).

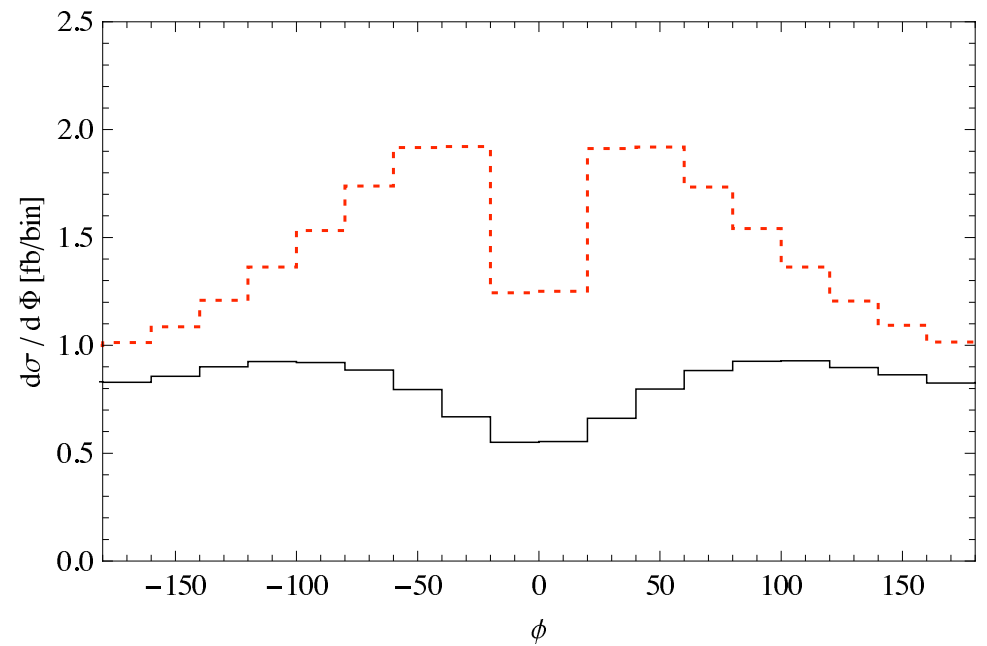
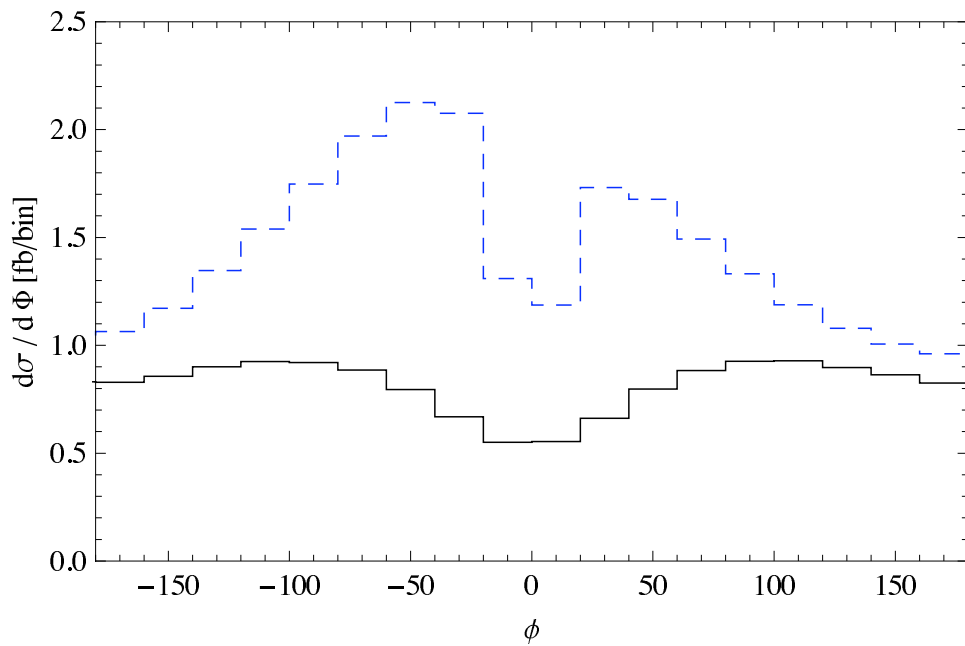
<sup>††</sup>Butterworth, Davison, Rubin, Salam (2008).

<sup>‡‡</sup>Plehn, Rainwater, Zeppenfeld (2002).

## Angular asymmetries at the LHC: $\text{Re}(\tilde{b})$ , $\text{Im}(\tilde{b})$

$$\phi_{ll} \equiv \text{sgn}((\vec{\ell}^+ - \vec{\ell}^-) \cdot \hat{z}) \sin^{-1}(\hat{\ell}^+ \times \hat{\ell}^- \cdot \hat{z}),$$

$$A_{\phi_{ll}} \equiv \frac{\sigma_{\phi_{ll} < 0} - \sigma_{\phi_{ll} > 0}}{\sigma_{\phi_{ll} < 0} + \sigma_{\phi_{ll} > 0}}$$

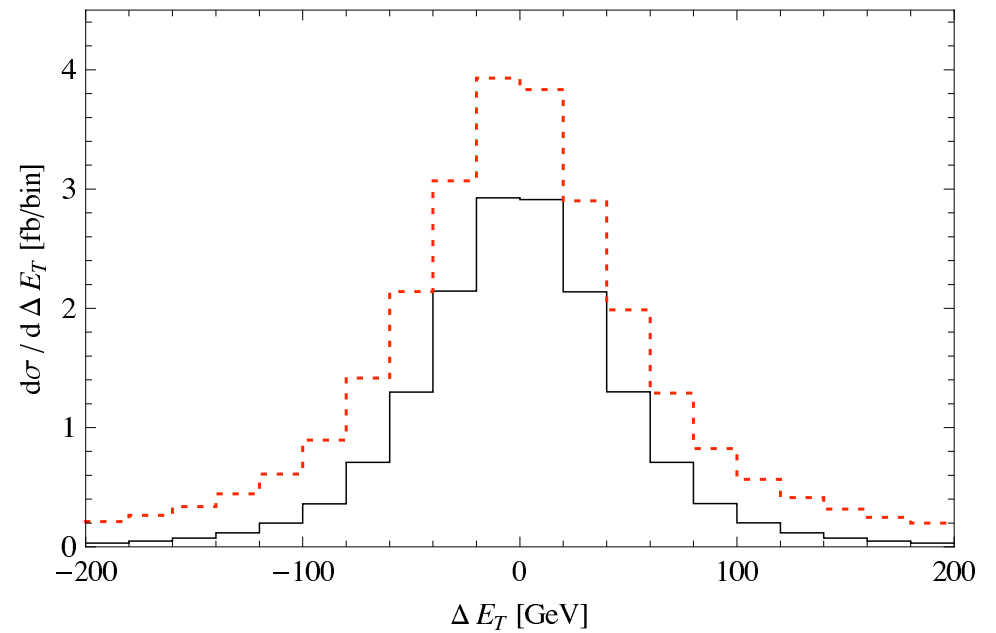
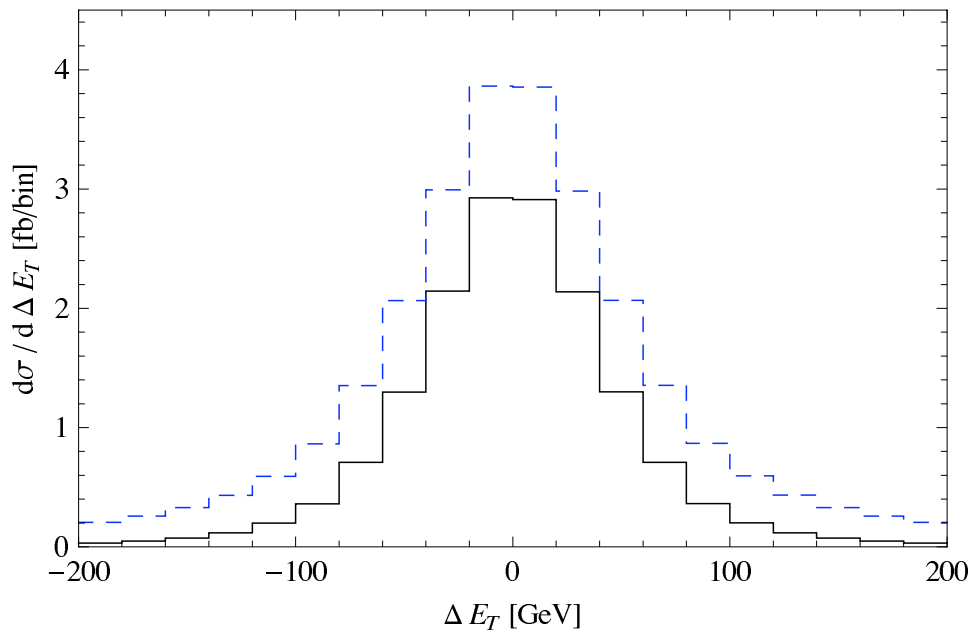


$$|\tilde{b}| = 0.25.$$

# Energy asymmetries at the LHC: $\text{Re}(\tilde{b})$ , $\text{Im}(\tilde{b})$

$$\Delta E_T \equiv E_T^+ - E_T^-$$

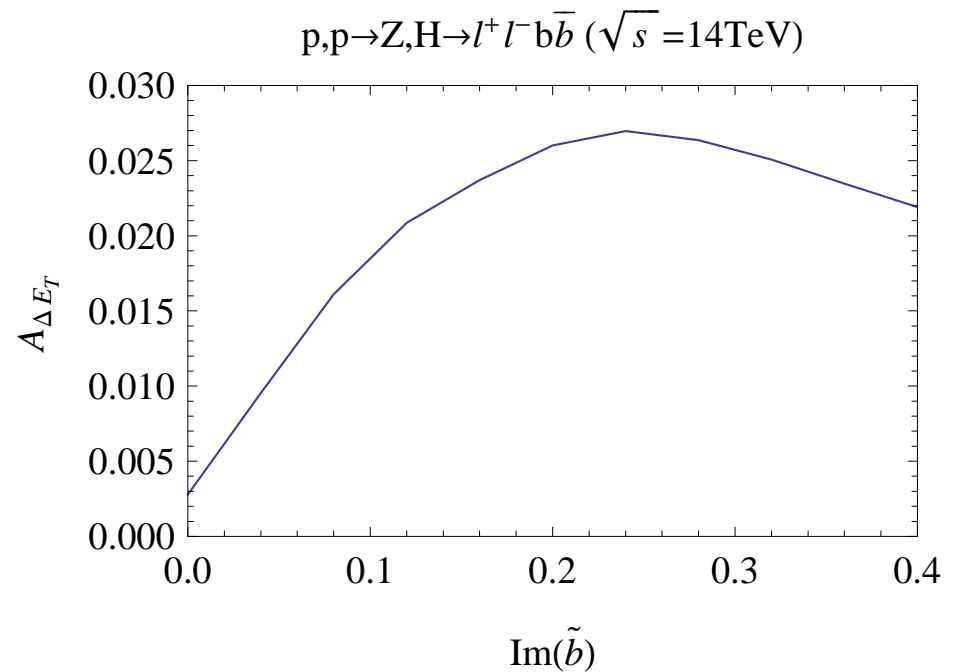
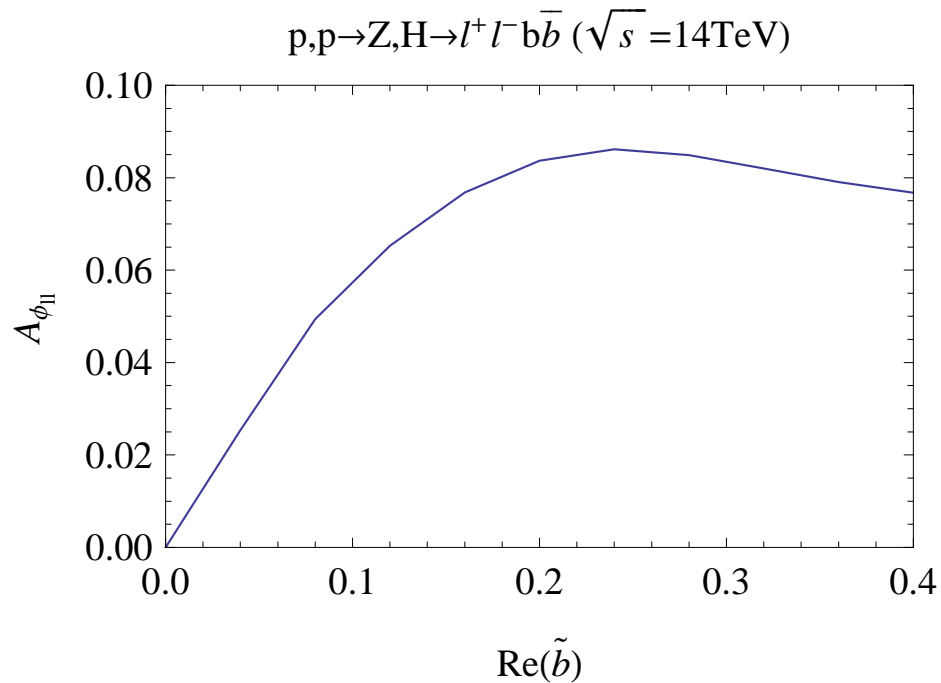
$$A_{\Delta E_T} \equiv \frac{\sigma_{\Delta E_T < 0} - \sigma_{\Delta E_T > 0}}{\sigma_{\Delta E_T < 0} + \sigma_{\Delta E_T > 0}}.$$



$$|\tilde{b}| = 0.25.$$



Possible size of asymmetries:  $\text{Re}(\tilde{b})$ ,  $\text{Im}(\tilde{b})$



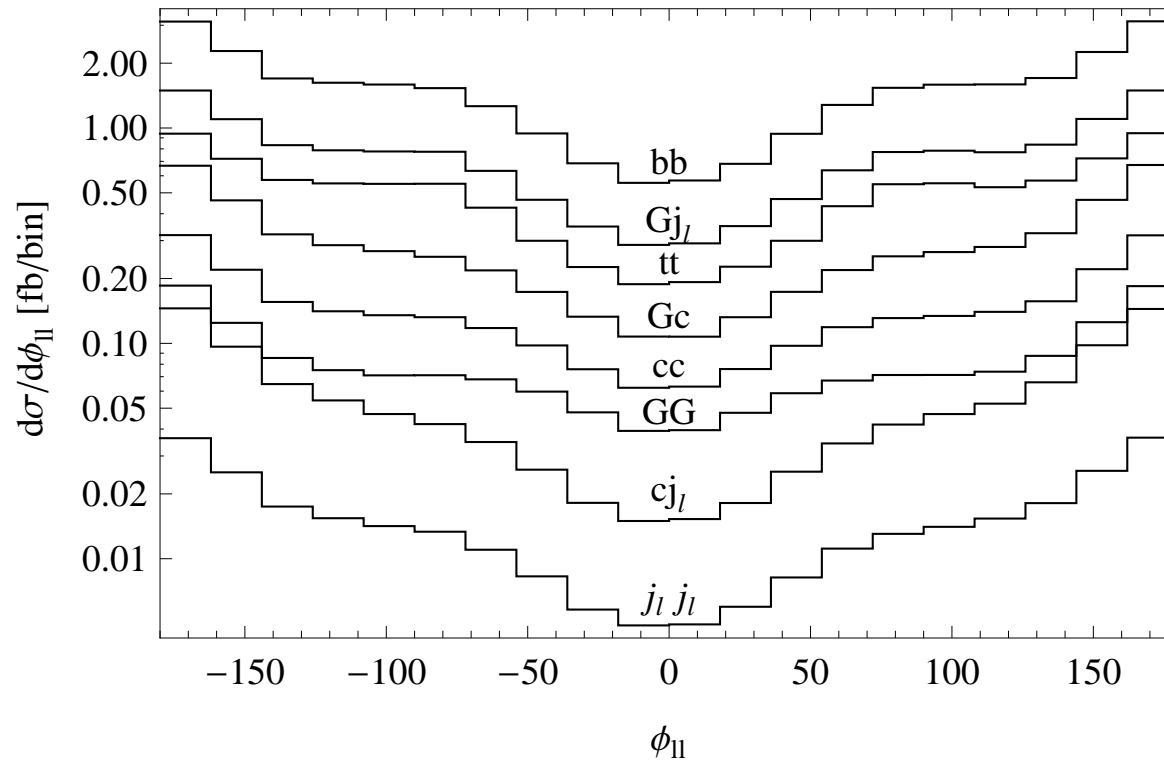
Around  $|\tilde{b}| = 0.25$ , linear approximation becomes a concern.

# $e^+e^- b\bar{b}$ backgrounds:

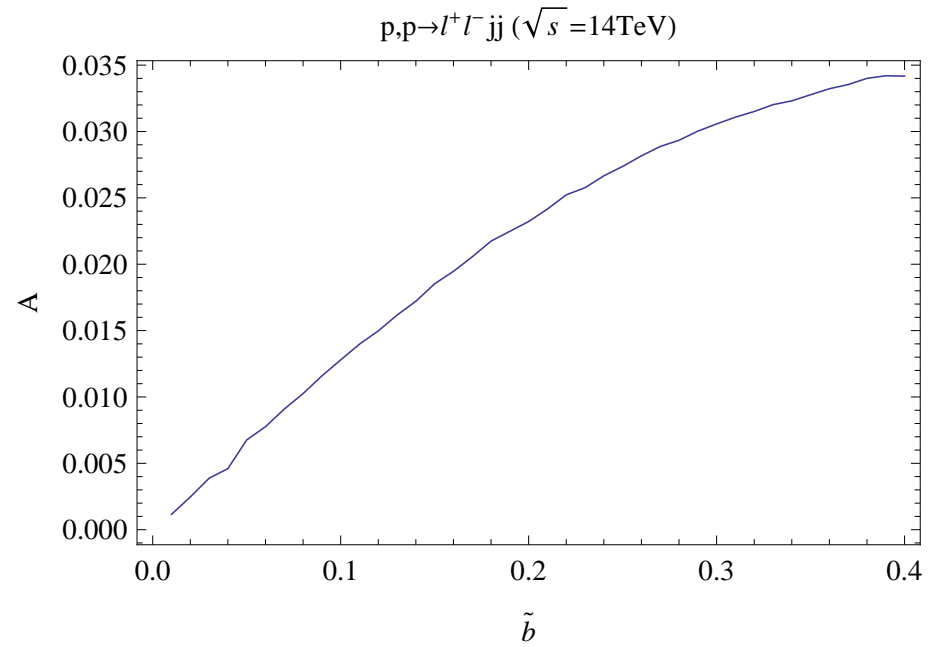
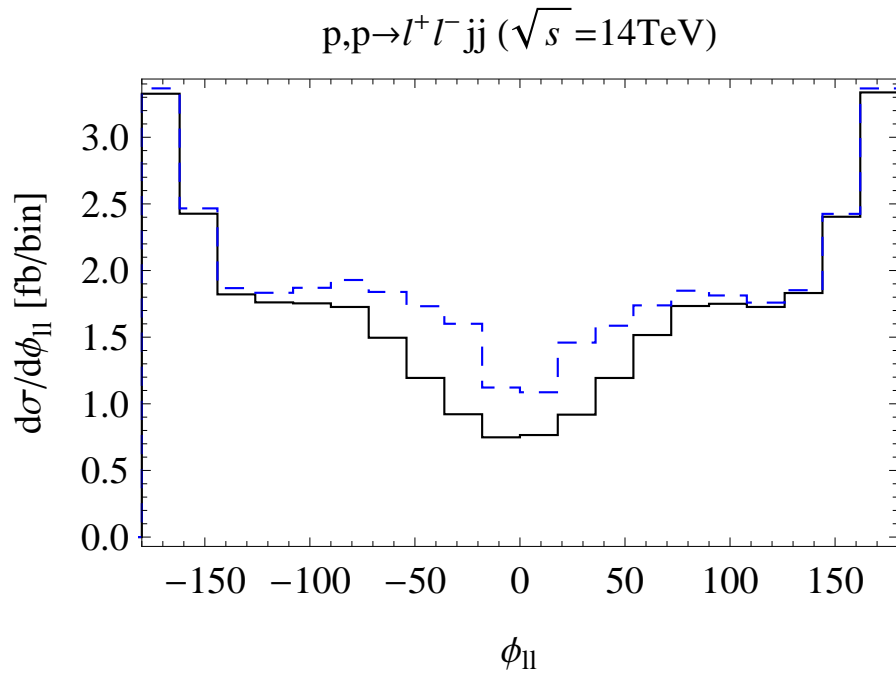
Jet tagging/mistagging rates (CMS TDR Figure 12.30):

b-quark jet	0.4
c-quark jet	0.03
Gluon jet	0.006
light jet	0.001

$p,p \rightarrow l^+l^- jj$  ( $\sqrt{s} = 14\text{TeV}$ )



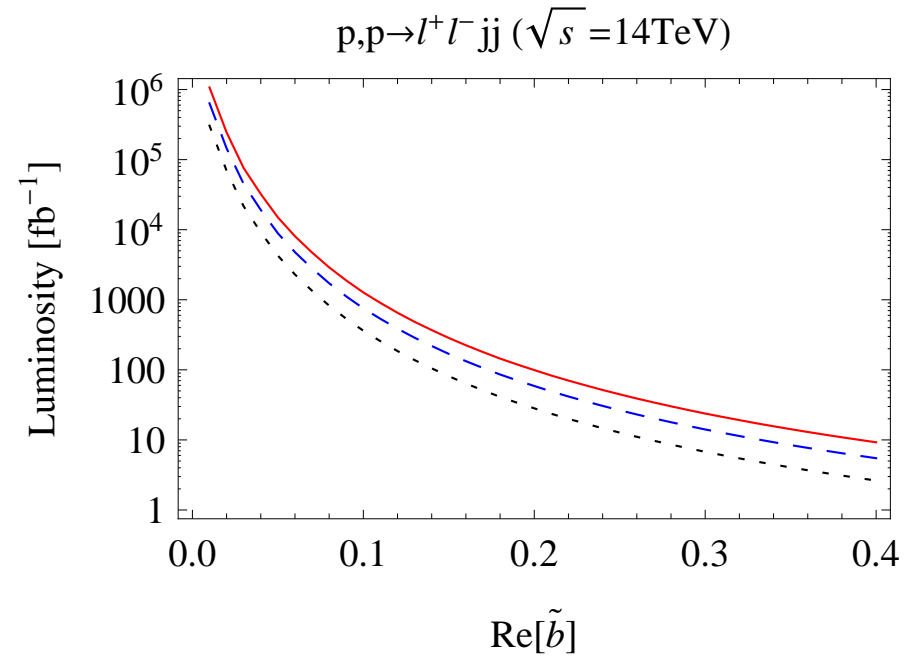
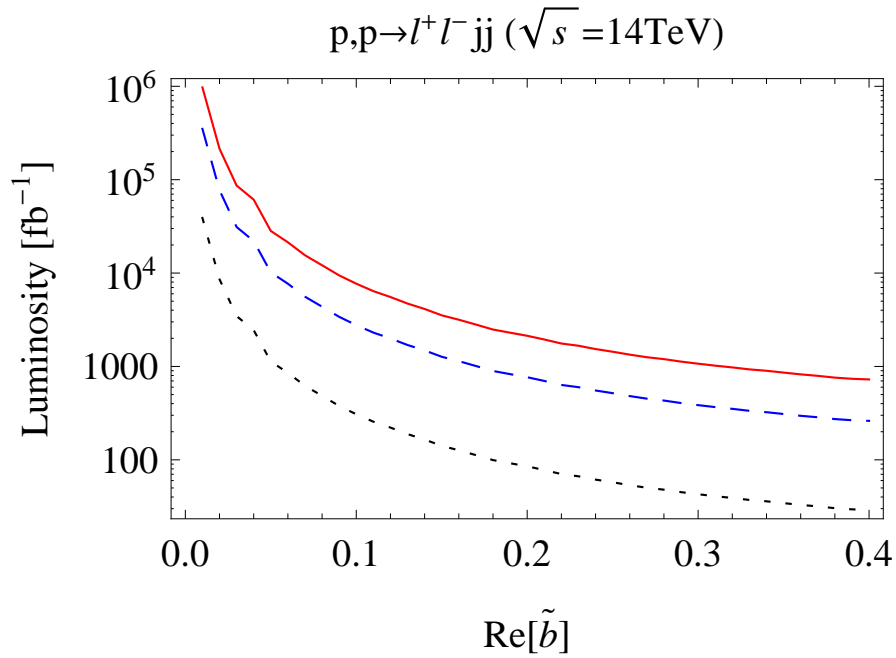
Including the backgrounds:



$$|\tilde{b}| = 0.25.$$

# Integrated luminosity needed for 1, 3, 5 $\sigma$

$$pp \rightarrow ZH \rightarrow \ell^+ \ell^- b\bar{b}$$



$$A \pm \sqrt{N}$$

Hypothesis test: Log likelihood  $LL(\tilde{b} \text{ vs } 0)$

## Concluding Remarks

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- Many possible application ...



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- As an example,  $pp \rightarrow ZH \rightarrow \ell^+\ell^- b\bar{b}$  has sensitivity to  $ZZH$  coupling.
- Many possible application ...

LHC to open up new physics era,  
and find “most wanted” new CP violation sources.